



On Generalized H-Birecurrent Finsler Space

Research Article

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Abstract: In this paper, we introduced a Finsler space for which the h-curvature tensor H_{jkh}^i (curvature tensor of Berwald) satisfies the condition

$$\mathcal{B}_m \mathcal{B}_n H_{jkh}^i = a_{mn} H_{jkh}^i + b_{mn} (\delta_j^i g_{kh} - \delta_k^i g_{jh}) - 2 y^r \mu_n \mathcal{B}_r (\delta_j^i C_{khm} - \delta_k^i C_{jhm}), H_{jkh}^i \neq 0$$

C_{jhm} is (h) hv-torsion tensor, where $\mathcal{B}_m \mathcal{B}_n$ is Berwald's covariant differential operator of the second order with respect to x^n and x^m , successively, \mathcal{B}_r is Berwald's covariant differential operator of the first order with respect to x^r , a_{mn} and b_{mn} are non-zero covariant tensors field of second order and μ_n is non-zero covariant vector field. We called this space a generalized H-birecurrent space. The aim of this paper is to develop some properties of a generalized H-birecurrent space by obtaining Berwald's covariant derivative of the second order for the (h)v-torsion tensor H_{jkh}^i and the deviation tensor H_j^i . The non-vanishing of Ricci tensor H_{kh} , the curvature vector H_k and the curvature scalar H are investigated. Different results regarding the covariant tensors field a_{mn} and b_{mn} have been established. Some conditions have been pointed out which reduce a generalized H-birecurrent space F_n ($n > 2$) into Landsberg space. We obtained an identity for a generalized H-birecurrent space. The conditions which reduce a generalized H-birecurrent space F_n ($n > 2$) in to a space of curvature scalar are given.

Keywords: Finsler space, Generalized H-birecurrent Finsler space, Ricci tensor, Landsberg space, Finsler space of scalar curvature.

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1. Introduction

A Finsler space of recurrent curvature was introduced and studied by P.N. Pandey [5, 6], P.N. Pandey and V.J. Dwivedi [7], R. Verma [9], S. Dikshit [10], F.Y.A. Qasem [1], N.S.H. Hussien [4], N.L. Youssef and A. Soleiman [3] and others. P.N. Pandey, S.S. Saxena and A. Goswami [6] introduced and studied a generalized H-recurrent Finsler space. Let F_n be an n-dimensional Finsler spaces equipped with the metric function F satisfies the requisite conditions [2].

Let the components of the corresponding metric tensor and Berwald's connection coefficients be denoted by g_{ij} and G_{jk}^i respectively. These are positively homogeneous of degree zero in the directional arguments. Due to their homogeneity in the directional arguments, we have ²

$$a) C_{ijk} y^i = C_{kij} y^i = C_{jki} y^i = 0 \text{ and } b) G_{jkh}^i y^j = G_{hjk}^i y^j = G_{kjh}^i y^j = 0, \quad (1)$$

where $C_{ijk} = \dot{\partial}_k g_{ij}$ and $G_{jkh}^i = \dot{\partial}_h G_{jk}^i$ are the components of tensors, they are symmetric in the their lower indices and $\dot{\partial}_h = \frac{\partial}{\partial y^h}$.

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¹ The indices i, j, k, \dots assume positive integral values from 1 to n.

² Unless stated otherwise. Henceforth all geometric object are assumed to be functions of line-element.

The relations between the metric function F , the components of the metric tensor and the vector y_i are given by

$$a) y_i = g_{ij}y^j, \quad b) g_{ij} = \dot{\partial}_i y_j = \dot{\partial}_j y_i, \quad \text{and} \quad c) y_i y^i = F^2. \quad (2)$$

The unit vector l^i in the direction of the directional argument is given by

$$a) l^i := \frac{y^i}{F} \quad (3)$$

and the associate vector of l_i is defined by

$$b) l_i := g_{ij}l^j = \dot{\partial}_i F = \frac{y_i}{F}.$$

Berwald covariant derivative of an arbitrary tensor field T_j^i with respect to x^h is given by

$$\mathcal{B}_h T_j^i := \partial_h T_j^i - \left(\dot{\partial}_r T_j^i \right) G_h^r + T_j^r G_{rh}^i - T_r^i G_{jh}^r. \quad (4)$$

Berwald covariant derivative of the vectors y^i and y_i with respect to x^k vanish identically i.e.

$$a) \mathcal{B}_k y^i = 0 \quad \text{and} \quad b) \mathcal{B}_k y_i = 0. \quad (5)$$

Berwald's covariant derivative of the metric tensor g_{ij} does not vanish in general ($\mathcal{B}_k g_{ij} \neq 0$) and is given by

$$\mathcal{B}_k g_{ij} = -2C_{ijk|r} y^r = -2y^h \mathcal{B}_h C_{ijk}, \quad (6)$$

where $|r$ is h-covariant derivative with respect to x^r (Cartan's second kind covariant differentiation). The processes of Berwald's covariant differentiation with respect to x^h and the partial differentiation with respect to directional argument y^k commute according to

$$\left(\dot{\partial}_k \mathcal{B}_h - \mathcal{B}_h \dot{\partial}_k \right) T_j^i = T_j^r G_{khr}^i - T_r^i G_{khr}^r \quad (7)$$

for an arbitrary tensor field T_j^i . The second covariant derivative of an arbitrary tensor field T_j^i with respect to x^h and x^k in the sense of Berwald may written as

$$\mathcal{B}_k \mathcal{B}_h T_j^i = \partial_k \mathcal{B}_h T_j^i - \left(\partial_s \mathcal{B}_h T_j^i \right) G_k^s + \left(\mathcal{B}_h T_j^r \right) G_{rk}^i - \left(\mathcal{B}_h T_r^i \right) G_{ik}^r - \left(\mathcal{B}_r T_j^i \right) G_{hk}^r. \quad (8)$$

The commutation formula for Berwald's curvature differentiation as follows:

$$\mathcal{B}_h \mathcal{B}_k T_j^i - \mathcal{B}_k \mathcal{B}_h T_j^i = T_j^r H_{hkr}^i - T_r^i H_{hkj}^r - \left(\dot{\partial}_r T_j^i \right) H_{hk}^r \quad (9)$$

where H_{jkh}^i ³ defined by

$$H_{jkh}^i = 2 \left\{ \partial_{[j} G_{k]h}^i + G_{rh[j}^i G_{k]h}^r + G_{r[j}^i G_{k]h}^r \right\} \quad (10)$$

are components of Berwald curvature tensor and

$$a) H_{jk}^i = H_{jkh}^i y^h \quad \text{and} \quad b) H_{jkh}^i = \dot{\partial}_h H_{jk}^i. \quad (11)$$

³ In Rund's book, H_{jkh}^i defined here, is denoted by H_{hjk}^i . This difference must be noted. The square brackets denote the skew-symmetric part of the tensor with respect to the indices enclosed therein.

It is clear from the definition that Berwald curvature tensor H_{jkh}^i is skew-symmetric in its first two lower indices and positively homogeneous of degree zero in the directional arguments y^i . Berwald's deviation tensor H_j^i is defined as

$$H_j^i = H_{jk}^i y^k. \quad (12)$$

The contraction of the indices i and j in H_{jkh}^i , H_{jk}^i and H_j^i gives

$$a) H_{kh} = H_{ikh}^i, \quad b) H_k = H_{ik}^i, \quad c) H = \frac{1}{n-1} H_i^i \quad \text{and} \quad d) H_k y^k = (n-1)H. \quad (13)$$

The necessary and sufficient condition for a Finsler space F_n ($n > 2$) to be Finsler space of scalar curvature is given by

$$H_h^i = F^2 R(\delta_h^i - l^i l_h). \quad (14)$$

2. Generalized H-Birecurrent Finsler Space

A Finsler space whose Berwald curvature tensor H_{jkh}^i satisfies the condition

$$\mathcal{B}_n H_{jkh}^i = \lambda_n H_{jkh}^i + \mu_n (\delta_j^i g_{kh} - \delta_k^i g_{jh}), \quad H_{jkh}^i \neq 0, \quad (A)$$

where \mathcal{B}_n is Berwald's covariant differential operator, λ_l and μ_n are non-zero covariant vector fields, this space introduced by P.N. Pandey, S. Saxena and A. Goswami [8], they called it as a generalized H-recurrent Finsler space. Now, taking the covariant derivative for the condition (A) with respect to x^m in the sense of Berwald and in view of the condition (A) and by using (5b), we get

$$\begin{aligned} \mathcal{B}_m \mathcal{B}_n H_{jkh}^i &= (\mathcal{B}_m \lambda_n) H_{jkh}^i + \lambda_n \left\{ \lambda_m H_{jkh}^i + \mu_m (\delta_j^i g_{kh} - \delta_k^i g_{jh}) \right\} \\ &+ (\mathcal{B}_m \mu_n) (\delta_j^i g_{kh} - \delta_k^i g_{jh}) + \mu_n \left\{ \delta_j^i (-2y^r \mathcal{B}_r C_{k h m}) - \delta_k^i (-2y^r \mathcal{B}_r C_{j h m}) \right\} \\ &= (\mathcal{B}_m \lambda_n + \lambda_n \lambda_m) H_{jkh}^i + (\lambda_n \mu_m + \mathcal{B}_m \mu_n) (\delta_j^i g_{kh} - \delta_k^i g_{jh}) - 2\mu_n y^r \mathcal{B}_r (\delta_j^i C_{k h m} - \delta_k^i C_{j h m}) \end{aligned}$$

which can be written as

$$\mathcal{B}_m \mathcal{B}_n H_{jkh}^i = a_{mn} H_{jkh}^i + b_{mn} (\delta_j^i g_{kh} - \delta_k^i g_{jh}) - 2\mu_n y^r \mathcal{B}_r (\delta_j^i C_{k h m} - \delta_k^i C_{j h m}), \quad H_{jkh}^i \neq 0, \quad (15)$$

where $a_{mn} = \mathcal{B}_m \lambda_n + \lambda_n \lambda_m$ and $b_{mn} = \lambda_n \mu_m + \mathcal{B}_m \mu_n$ are non-zero covariant tensors field of second order are non-null covariant tensors field of second order.

Definition 2.1. A Finsler space F_n whose Berwald curvature tensor H_{jkh}^i satisfies the condition (15), where a_{mn} and b_{mn} are non-zero covariant tensors field of second order. We shall call such Finsler space as a generalized H-birecurrent Finsler space and briefly denoted by $GH - BRF_n$.

Let us consider a $GH - BRF_n$ which is characterized by the condition (15). Transvecting the condition (15) by y^h , using (11a), (2a), (5a) and (1a), we get

$$\mathcal{B}_m \mathcal{B}_n H_{jk}^i = a_{mn} H_{jk}^i + b_{mn} (\delta_j^i y_k - \delta_k^i y_j). \quad (16)$$

Transvecting the condition (16) by y^k , using (12), (5a) and (2c), we get

$$\mathcal{B}_m \mathcal{B}_n H_j^i = a_{mn} H_j^i + b_{mn} (\delta_j^i F^2 - y_j y^i). \quad (17)$$

Thus, we conclude

Theorem 2.2. In $GH - BRF_n$, Berwald covariant derivative of the second order for the $h(v)$ -torsion tensor H_{kh}^i and the deviation tensor H_h^i given by the conditions (16) and (17), respectively.

Contracting the indices i and j in the conditions (15), (16) and (17), using (13a), (13b), (13c) and (2c), we get

$$\mathcal{B}_m \mathcal{B}_n H_{kh} = a_{mn} H_{kh} + (n-1) b_{mn} g_{kh} - 2(n-1) \mu_n y^r \mathcal{B}_r C_{khm}, \quad (18)$$

$$\mathcal{B}_m \mathcal{B}_n H_k = a_{mn} H_k + (n-1) b_{mn} y_k \quad (19)$$

and

$$\mathcal{B}_m \mathcal{B}_n H = a_{mn} H + b_{mn} F^2. \quad (20)$$

The conditions (18), (19) and (20) show that Ricci tensor H_{kh} , the curvature vector H_k and the curvature scalar H can't vanish because the vanishing of any one of these would imply $a_{mn} = 0$, $b_{mn} = 0$ and $\mu_n = 0$, that is a contradiction. Thus, we conclude

Theorem 2.3. In $GH - BRF_n$, Ricci tensor H_{kh} , the curvature vector H_k and the curvature scalar H are non-vanishing.

Let us consider a $GH - BRF_n$. Differentiating the condition (19) partially with respect to y^h , using $(\dot{\partial}_h H_k = H_{kh})$ and (2b), we get

$$\dot{\partial}_h \mathcal{B}_m (\mathcal{B}_n H_k) = (\dot{\partial}_h a_{mn}) H_k + a_{mn} H_{kh} + (n-1) (\dot{\partial}_h b_{mn}) y_k + (n-1) b_{mn} g_{kh}. \quad (21)$$

Using the commutation formula (7) for $(\mathcal{B}_n H_k)$ in (21), we get

$$\mathcal{B}_m \dot{\partial}_h (\mathcal{B}_n H_k) - (\mathcal{B}_r H_k) G_{hmn}^r - (\mathcal{B}_n H_r) G_{hmk}^r = (\dot{\partial}_h a_{mn}) H_k + a_{mn} H_{kh} + (n-1) (\dot{\partial}_h b_{mn}) y_k + (n-1) b_{mn} g_{kh}. \quad (22)$$

Again applying the commutation formula (7) for (H_k) in (22) and $(\dot{\partial}_h H_k = H_{kh})$, we get

$$\mathcal{B}_m \mathcal{B}_n H_{kh} - \mathcal{B}_m (H_r G_{hnk}^r) - (\mathcal{B}_r H_k) G_{hmn}^r - (\mathcal{B}_n H_r) G_{hmk}^r = (\dot{\partial}_h a_{mn}) H_k + a_{mn} H_{kh} + (n-1) (\dot{\partial}_h b_{mn}) y_k + (n-1) b_{mn} g_{kh}. \quad (23)$$

Using the condition (18) in (23), we get

$$-2(n-1) \mu_n y^r \mathcal{B}_r C_{khm} - \mathcal{B}_m (H_r G_{hnk}^r) - (\mathcal{B}_r H_k) G_{hmn}^r - (\mathcal{B}_n H_r) G_{hmk}^r = (\dot{\partial}_h a_{mn}) H_k + (n-1) (\dot{\partial}_h b_{mn}) y_k. \quad (24)$$

Transvecting (24) by y^k , using (5a), (1a), (1b), (13d) and (2c), we get

$$- (\mathcal{B}_r H) G_{hmn}^r = (\dot{\partial}_h a_{mn}) H + (\dot{\partial}_h b_{mn}) F^2. \quad (25)$$

If $(\mathcal{B}_r H) G_{hmn}^r = 0$, the equation (25) can be written as

$$\dot{\partial}_h b_{mn} = - \frac{\dot{\partial}_h a_{mn}}{F^2} H. \quad (26)$$

If the covariant tensor field a_{mn} is independent of the directional argument, the equation (26) shows that the covariant tensor field b_{mn} is also independent of the directional argument. Conversely, if the covariant tensor field b_{mn} is independent of the directional argument, we get $(\dot{\partial}_h a_{mn}) H = 0$. In view of Theorem 2.2, the condition $(\dot{\partial}_h a_{mn}) H = 0$ implies $\dot{\partial}_h a_{mn} = 0$, i.e. the covariant tensor field a_{mn} is also independent of the directional argument. Thus, we conclude

Theorem 2.4. *In $GH - BRF_n$, the covariant tensor field b_{mn} is independent of the directional argument if and only if the covariant tensor field a_{mn} is independent of the directional argument provided $(\mathcal{B}_r H) G_{hmn}^r = 0$.*

Transvecting (25) by y^m and using (1b), we get

$$\dot{\partial}_h b_n - b_{hn} = -\frac{(\dot{\partial}_h a_n - a_{hn})}{F^2} H, \quad (27)$$

where $a_{mn} y^m = a_n$ and $b_{mn} y^m = b_n$. Since the covariant vector field a_n is not independent of the directional argument, the equation (27) shows that the covariant vector field b_n is also not independent of the directional argument. Conversely, if the covariant vector field b_n is not independent of the directional argument, we have $(\dot{\partial}_h a_n - a_{hn}) H = 0$. In view of Theorem 2.2, the condition $(\dot{\partial}_h a_n - a_{hn}) H = 0$ implies $\dot{\partial}_h a_n = a_{hn}$, i.e. the covariant vector field a_n also is not independent of the directional argument. Thus, we conclude

Theorem 2.5. *In $GH - BRF_n$, the covariant vector field a_n is not independent of the directional argument if and only if the covariant vector field b_n is not independent of the directional argument.*

Suppose that the covariant tensor field a_{mn} is not independent of the directional argument, then by using (26) in (24), we have

$$-2(n-1)\mu_n y^r \mathcal{B}_r C_{khn} - \mathcal{B}_m (H_r G_{hnk}^r) - (\mathcal{B}_r H_k) G_{hmn}^r - (\mathcal{B}_n H_r) G_{hmk}^r = (\dot{\partial}_h a_{mn}) \left\{ H_k - \frac{(n-1)H}{F^2} y_k \right\}. \quad (28)$$

Transvecting (28) by y^k , using (5a), (1a), (1b), (13d) and (2c), we get

$$(\mathcal{B}_r H) G_{hmn}^r = 0.$$

Thus, we have

Theorem 2.6. *In $GH - BRF_n$, we have the identity $(\mathcal{B}_r H) G_{hmn}^r = 0$.*

Transvecting (28) by y^m and using (1a) and (1b), we get

$$-\mathcal{B}_m (H_r G_{hnk}^r) y^m = (\dot{\partial}_h a_n - a_{hn}) \left\{ H_k - \frac{(n-1)H}{F^2} y_k \right\}, \quad (29)$$

where $a_{mn} y^m = a_n$. If $\mathcal{B}_m (H_r G_{hnk}^r) y^m = 0$, the equation (29) implies at least one of the following conditions

$$a) \ a_{hn} = \dot{\partial}_h a_n \quad \text{or} \quad b) \ H_k = \frac{(n-1)H}{F^2} y_k. \quad (30)$$

Thus, we conclude

Theorem 2.7. *In $GH - BRF_n$, for which the covariant tensor field a_{hn} is not independent of the directional argument at least one of the conditions (30a) or (30b) holds provided $\mathcal{B}_m (H_r G_{hnk}^r) y^m = 0$.*

Suppose that (30b) holds. Then (28) implies

$$-2\mu_n y^r \mathcal{B}_r C_{khn} - \mathcal{B}_m \left(\frac{H}{F^2} y_r G_{hnk}^r \right) - (\mathcal{B}_r \frac{H}{F^2} y_k) G_{hmn}^r - (\mathcal{B}_n \frac{H}{F^2} y_r) G_{hmk}^r = 0. \quad (31)$$

Transvecting (31) by y^m , using (5a), (1a) and (1b), we get

$$\mathcal{B}_m \left(\frac{H}{F^2} y_r G_{hnk}^r \right) y^m = 0$$

which can be written as

$$F^{-2} \{(\mathcal{B}_m H) y_r G_{hnk}^r + H \mathcal{B}_m (y_r G_{hnk}^r)\} y^m = 0.$$

If $H \mathcal{B}_m (y_r G_{hnk}^r) y^m = 0$, the above equation implies $(\mathcal{B}_m H) y^m y_r G_{hnk}^r = 0$. Since $(\mathcal{B}_m H) y^m \neq 0$, so $y_r G_{hnk}^r = 0$, therefore the space is a Landsberg space. Thus, we conclude

Theorem 2.8. *The GH – BRF_n is a Landsberg space if the condition (30b) holds provided $\mathcal{B}_m (H_r G_{hnk}^r) y^m = 0$ and $(\mathcal{B}_m H) y^m y_r G_{hnk}^r = 0$.*

If the covariant tensor field $a_{hn} \neq \dot{\partial}_h a_n$, in view of Theorem 2.4 (30b) holds good. In view of this fact, we may rewrite Theorem 2.7 in the following form

Theorem 2.9. *The GH – BRF_n is necessarily a Landsberg space provided $a_{hn} \neq \dot{\partial}_h a_n$, $\mathcal{B}_m (H_r G_{hnk}^r) y^m = 0$ and $(\mathcal{B}_m H) y^m y_r G_{hnk}^r = 0$.*

Differentiating the condition (16) partially with respect to y^h , using (11b) and (2b), we get

$$\dot{\partial}_h \mathcal{B}_m (\mathcal{B}_n H_{jk}^i) = (\dot{\partial}_h a_{mn}) H_{jk}^i + a_{mn} H_{jkh}^i + (\dot{\partial}_h b_{mn}) (\delta_j^i y_k - \delta_k^i y_j) + b_{mn} (\delta_j^i g_{kh} - \delta_k^i g_{jh}). \quad (32)$$

Using the commutation formula (7) for $(\mathcal{B}_n H_{jk}^i)$ in (32), we get

$$\begin{aligned} \mathcal{B}_m (\dot{\partial}_h \mathcal{B}_n H_{jk}^i) - (\mathcal{B}_r H_{jk}^i) G_{hmn}^r + (\mathcal{B}_n H_{jk}^r) G_{hmr}^i - (\mathcal{B}_n H_{rk}^i) G_{hmj}^r - (\mathcal{B}_n H_{jr}^i) G_{hmk}^r \\ = (\dot{\partial}_h a_{mn}) H_{jk}^i + a_{mn} H_{jkh}^i + (\dot{\partial}_h b_{mn}) (\delta_j^i y_k - \delta_k^i y_j) + b_{mn} (\delta_j^i g_{kh} - \delta_k^i g_{jh}). \end{aligned} \quad (33)$$

Again using the commutation formula (7) for (H_{jk}^i) in (33) and using (11b), we get

$$\begin{aligned} \mathcal{B}_m \mathcal{B}_n H_{jkh}^i + (\mathcal{B}_m H_{jk}^r) G_{hnr}^i + H_{jk}^r (\mathcal{B}_m G_{hnr}^i) - (\mathcal{B}_m H_{rk}^i) G_{hnr}^r - H_{rk}^i (\mathcal{B}_m G_{hnr}^r) - (\mathcal{B}_m H_{jr}^i) G_{hnk}^r - H_{jr}^i (\mathcal{B}_m G_{hnk}^r) \\ - (\mathcal{B}_r H_{jk}^i) G_{hmn}^r + (\mathcal{B}_n H_{jk}^r) G_{hmr}^i - (\mathcal{B}_n H_{jr}^i) G_{hmk}^r = (\dot{\partial}_h a_{mn}) H_{jk}^i + a_{mn} H_{jkh}^i - (\mathcal{B}_n H_{rk}^i) G_{hmj}^r \\ + (\dot{\partial}_h b_{mn}) (\delta_j^i y_k - \delta_k^i y_j) + b_{mn} (\delta_j^i g_{kh} - \delta_k^i g_{jh}). \end{aligned}$$

By using the condition (15), the above equation can be written as

$$\begin{aligned} (\mathcal{B}_m H_{jk}^r) G_{hnr}^i + H_{jk}^r (\mathcal{B}_m G_{hnr}^i) - (\mathcal{B}_m H_{rk}^i) G_{hnr}^r - H_{rk}^i (\mathcal{B}_m G_{hnr}^r) - (\mathcal{B}_m H_{jr}^i) G_{hnk}^r - H_{jr}^i (\mathcal{B}_m G_{hnk}^r) - (\mathcal{B}_r H_{jk}^i) G_{hmn}^r \\ + (\mathcal{B}_n H_{jk}^r) G_{hmr}^i - (\mathcal{B}_n H_{rk}^i) G_{hmj}^r - (\mathcal{B}_n H_{jr}^i) G_{hmk}^r \\ = (\dot{\partial}_h a_{mn}) H_{jk}^i + (\dot{\partial}_h b_{mn}) (\delta_j^i y_k - \delta_k^i y_j) + 2 \mu_n y^r \mathcal{B}_r (\delta_j^i C_{khn} - \delta_k^i C_{jhm}). \end{aligned} \quad (34)$$

Transvecting (34) by y_i , using the identity $(y_i H_{jk}^i = 0)$ which established by [7] and (5b), we get

$$(\mathcal{B}_m H_{jk}^r) y_i G_{hnr}^i + H_{jk}^r \{ \mathcal{B}_m (y_i G_{hnr}^i) \} + (\mathcal{B}_n H_{jk}^r) y_i G_{hmr}^i = 2 \mu_n y^r \mathcal{B}_r (y_j C_{khn} - y_k C_{jhm}). \quad (35)$$

Transvecting (35) by y^m , using (1b), (5a) and (1a), we get

$$\left[(\mathcal{B}_m H_{jk}^r) y_i G_{hnr}^i + H_{jk}^r \{ \mathcal{B}_m (y_i G_{hnr}^i) \} \right] y^m = 0$$

which can be written as

$$\left[\mathcal{B}_m (H_{jk}^r y_i G_{hnr}^i) \right] y^m = 0.$$

Thus, we conclude

Theorem 2.10. In $GH - BRF_n$, we have the identity $[\mathcal{B}_m (H_{jk}^r y_i G_{hnr}^i)] y^m = 0$.

Transvecting (34) by y^k , using (12), (1b), (5a), (2c) and (1a), we get

$$\begin{aligned} & (\mathcal{B}_m H_j^r) G_{hnr}^i + H_j^r (\mathcal{B}_m G_{hnr}^i) - (\mathcal{B}_m H_r^i) G_{hnr}^r - H_r^i (\mathcal{B}_m G_{hnr}^r) - (\mathcal{B}_r H_j^i) G_{hmn}^r + (\mathcal{B}_n H_j^r) G_{hmr}^i - (\mathcal{B}_n H_r^i) G_{hmj}^r \\ & = (\dot{\partial}_h a_{mn}) H_j^i + (\dot{\partial}_h b_{mn}) (\delta_j^i F^2 - y^i y_j) + 2 \mu_n y^r \mathcal{B}_r (y^i C_{jhm}). \end{aligned} \tag{36}$$

Substituting the value of $(\dot{\partial}_h b_{mn})$ from (26) in (36), using (3a) and (3b), we get

$$\begin{aligned} & \mathcal{B}_m (H_j^r G_{hnr}^i) - \mathcal{B}_m (H_r^i G_{hnr}^r) - (\mathcal{B}_r H_j^i) G_{hmn}^r + (\mathcal{B}_n H_j^r) G_{hmr}^i - (\mathcal{B}_n H_r^i) G_{hmj}^r \\ & = (\dot{\partial}_h a_{mn}) \{H_j^i - H (\delta_j^i - l^i l_j)\} + 2 \mu_n y^r \mathcal{B}_r (y^i C_{jhm}). \end{aligned}$$

$$\text{If } \mathcal{B}_m (H_j^r G_{hnr}^i) - \mathcal{B}_m (H_r^i G_{hnr}^r) \mathcal{B}_m - (\mathcal{B}_r H_j^i) G_{hmn}^r + (\mathcal{B}_n H_j^r) G_{hmr}^i - (\mathcal{B}_n H_r^i) G_{hmi}^r - 2\mu_n y^r \mathcal{B}_r (y^i C_{jhm}) = 0. \tag{37}$$

Then, we have at least one of the following conditions

$$a) \dot{\partial}_h a_{hn} = 0 \text{ or } b) H_j^i = H (\delta_j^i - l^i l_j). \tag{38}$$

Putting $H = F^2 R$, (2.24b) may be written as

$$H_j^i = F^2 R (\delta_j^i - l^i l_j).$$

Therefore, the space is a Finsler space of scalar curvature. Thus, we conclude

Theorem 2.11. In $GH - BRF_n$, for $n > 2$ admitting the condition (37) is a Finsler space of scalar curvature provided $R \neq 0$ and the covariant tensor field a_{mn} is not independent of the directional argument.

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