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# Fractional Complex Variables: Strong Local Fractional Complex Derivatives (LFCDs) of Non-Integer Rational Order

**Review Article** 

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**Abstract:** Fractional complex variables focus on the fractional or non-integer order differential calculus of a complex variable. In fractional calculus, locality can narrow down pieces of a function where there may be better behavior in order to model in an analytic sense, as well as obtain more meaningful physical and/or geometric information. That's where we introduce the concepts of Strong Local Fractional Complex Derivatives or LFCDs. Strong LFCDs can "maximize" the opportunity that the piece of the function in a localized or local enough area is "well-behaved" (enough). We prove a theorem that shows where Strong LFCDs exist. Applications include index of stability in Complex or Real Fractional Advection Dispersion Equation (FADE).

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## 1. Introduction to Fractional Complex Derivatives

We start by looking at derivatives of fractional or non-integer order s.t  $\alpha \in \mathbb{Q}$ 

$$\frac{d^{\alpha}y}{dx^{\alpha}}$$
 or  $\frac{d^{\alpha}f(x)}{dx^{\alpha}}, \frac{d^{\frac{1}{2}}f(x)}{dx^{\frac{1}{2}}}$ 

**Definition 1.1.** If for a function  $f:[0,1] \to \mathbb{R}$ , then  $\exists$  the limit defining the derivative, where  $\alpha$  is  $0 < \alpha < 1$ 

$$D^{\alpha}f(y) = \lim_{x \to y} \frac{d^{\alpha}(f(x) - f(y))}{d(x - y)^{\alpha}}$$

This is the Local Fractional Derivative (LFD) form: If for a function  $f : [0,1] \to \mathbb{R}$ , then  $\exists$  a finite limit, where N is the largest integer for which Nth derivative of f(x) at y exists and is finite, then we say that the LFD of order  $\alpha : 1 \le \alpha \le N$  at x = y exists.

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$$D^{\alpha}f(y) = \lim_{x \to y} \left[ \frac{d^{\alpha} \left[ f(x) - \sum_{n=0}^{N} \frac{f^{(n)}(y)}{\Gamma(n+1)} f(x-y)^{n} \right]}{d[x-y]^{n}} \right]$$

Generally, fractional derivatives are not local in nature; however, if we localize some function f, we can use LDF to solve some physical models that integer order derivatives cannot really solve or explain. Now, recall W = f(z):

$$\begin{split} z &= x + iy, Re(z) = x, Im(z) = y \\ w &= u(x,y) + iv(x,y) = f(z) \\ \Rightarrow \lim_{z \to a + bi} f(z) &= \lim_{(x,y) \to (a,b)} u(x,y) + i \lim_{(x,y) \to (a,b)} v(x,y) \end{split}$$

where  $z \in \mathbb{C}, u, v \in \mathbb{R}$ 

 $\Rightarrow \text{Let } z \in R, \text{ then we have } z^{\alpha} \in R^{\alpha} \subseteq \mathbb{C}.$ And if  $z^{\alpha} \exists \Rightarrow w^{\alpha}.$ 

$$\Rightarrow w^{\alpha} = f(z) = u(x, y) + i^{\alpha}v(x, y)$$

Let  $f: F \to R^{\alpha}$  local function defined on a fractal set F of fractal dimension  $\alpha, 0 \leq \alpha \leq 1$ . If  $\forall \varepsilon > 0, \exists$  some  $\delta > 0$  s.t.

$$|f(z) - L| < \varepsilon^{\alpha} \Rightarrow 0 < |z - z_0| < \delta$$

The limit of f(z) as  $z \to z_0$  is L

$$\Rightarrow \lim_{z \to z_0} f(z) = L.$$

The function f(z) is said to be local fractional continuous at  $z_0$  if  $f(z_0)$  is defined, and

$$\lim_{z \to z_0} f(z) = f(z_0).$$

A function f(z) is deemed local fractional cont. on  $R^{\alpha}$  if it is local fractional continuous  $\forall$  point of its domain  $\mathbb{C}_{\alpha}(R)$ .

Let the complex function f(z) be defined in a neighborhood of a point  $z_0$ . The local fractional complex derivative of f(z) at  $z_0$  denoted by

$$D^{\alpha}f(z), \left.\frac{d^{\alpha}}{dz^{\alpha}}f(z)\right|_{z=z_{0}} \text{ or } f^{(\alpha)}(z_{0}),$$
$$= \lim_{z \to z_{0}} \frac{\Gamma(1+\alpha)[f(z) - f(z_{0})]}{(z-z_{0})^{\alpha}}, 0 < \alpha \leq 1$$
(1)

If this limit exists, then the function f(z) is said to be local fractional analytic at  $z_0$ . If this limit exists  $\forall z_0 \in \mathbb{R}^{\alpha}$ , then the function f(z) is deemed to be local fractional analytic in  $\mathbb{R}^{\alpha}$ .

If  $\exists$  a function

$$f(z) = u(x,y) + i^{\alpha}v(x,y) \tag{2}$$

The local fractional equations

$$\frac{\partial^{\alpha} u(x,y)}{\partial x^{\alpha}} - \frac{\partial^{\alpha} v(x,y)}{\partial y^{\alpha}} = 0$$
(3)

$$\frac{\partial^{\alpha} u(x,y)}{\partial y^{\alpha}} + \frac{\partial^{\alpha} v(x,y)}{\partial x^{\alpha}} = 0$$
(4)

are local fractional Cauchy-Riemann Equations.

# 2. Strong (or Weak) Local Fractional Complex Derivatives

**Theorem 2.1** (from Yang). Suppose that (2) is local fractional analytic in a region  $\mathbb{R}^{\alpha}$ . Then we have (3) and (4).

*Proof.* Local fractional C-R Equations are sufficient equations/conditions that f(z) be local functional analytic in region  $R^{\alpha} \Rightarrow R^{1}$ , where  $\alpha$  is 1. The local fractional partial equations

$$\frac{\partial^{2\alpha}u(x,y)}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha}u(x,y)}{\partial y^{2\alpha}} = 0$$
(5)

$$\frac{\partial^{2\alpha}v(x,y)}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha}v(x,y)}{\partial y^{2\alpha}} = 0 \tag{6}$$

are deemed local fractional Laplace Equations, denoted by

 $\nabla$ 

$$^{\alpha}u(x,y) = 0, \nabla^{\alpha}v(x,y) = 0 \ (\neg a,b)$$

$$\Rightarrow \nabla^{\alpha} = \frac{\partial^{2\alpha}}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha}}{\partial y^{2\alpha}}$$

$$\nabla^{\alpha} = \frac{\partial^{2\alpha}}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha}}{\partial y^{2\alpha}}$$
(7)

This is a local fractional Laplace operator. Suppose  $\nabla^{\alpha} u(x, y) = 0$ , then u(x, y) is a local fractional harmonic function in R. When local may not be good enough, there may be cases where physically, geometrically - a strong local fractional complex derivative or strong LFCD may be needed. Hence, we have maybe the following:

**Theorem 2.2** (Theorem Proposition (Strong LFCDs)). If  $\exists$  LFCDs of non-integer order  $\alpha \in \mathbb{Q}$  s.t.  $\alpha : 1 \leq \alpha \leq 2$ , then these are Strong LFCDs.

*Proof.* If a function  $f(z_0) \in \mathbb{C}$  domain is sufficiently smooth, and it meets Cauchy-Riemann conditions then  $D^{\alpha}(f(z))$  at least exists  $\forall \alpha \in \mathbb{Q}$  s.t.  $1 \leq \alpha < \infty$ . See previous Theorem.

Now, recall w = f(z) = u(x, y) + iv(x, y) and  $\exists$  partial derivatives for f with x and y so we have the following for  $f(z_0)$ :

$$f'(z_0) = \left(\frac{\partial u}{\partial x} + i\frac{\partial v}{\partial x}\right)_{(x_0, y_0)}$$
$$f'(z_0) = \left(-i\frac{\partial u}{\partial y} + \frac{\partial v}{\partial y}\right)_{(x_0, y_0)}$$
$$\Rightarrow \left(\frac{\partial u}{\partial x} + i\frac{\partial v}{\partial x}\right)_{(x_0, y_0)} = \left(-i\frac{\partial u}{\partial y} + \frac{\partial v}{\partial y}\right)_{(x_0, y_0)}$$
$$\Rightarrow \frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}, \ \frac{\partial v}{\partial x} = -\frac{\partial u}{\partial y}$$

These are the C-R Equations with  $\alpha = 1$ . By previous Theorem from Yang,  $\exists$  local fractional Laplace equations (5) and (6)

$$\frac{\partial^{2\alpha} u}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha} u}{\partial y^{2\alpha}} \ , \ \frac{\partial^{2\alpha} v}{\partial x^{2\alpha}} + \frac{\partial^{2\alpha} v}{\partial y^{2\alpha}}$$

with  $\alpha = 1$ 

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = \frac{\partial}{\partial x} \left( \frac{\partial v}{\partial y} \right) + \frac{\partial}{\partial y} \left( -\frac{\partial v}{\partial x} \right) = 0$$
$$\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} = \frac{\partial}{\partial x} \left( -\frac{\partial u}{\partial y} \right) + \frac{\partial}{\partial y} \left( \frac{\partial u}{\partial x} \right) = 0$$

This shows the existence of Laplace's Equation of order  $\alpha = 2$  is well established and defined for a local enough second order f(z) and we know C-R Equations of order 1 is also well-established and defined for a local first order f(z), then  $\Rightarrow$  we may have a strong or strong enough LCFDs which can exist for every  $\alpha : 1 \le \alpha \le 2 \ \forall \alpha \in \mathbb{Q}$  and  $z_0$  in  $f(z_0) \in \mathbb{C}$ .

**Remark 2.3.** On the contrary, LFCDs with order  $\alpha < 1$  and  $\alpha > 2$   $\forall \alpha \in \mathbb{Q}$  are not strong (enough) or are even weak. As a consequence of Theorem for strong LFCDs proposition.

For  $\alpha = 1$ ,  $\alpha = 2$ , the strong case seems evident. Using Sobolev Spaces and The Sharp Trace Theorem where  $\exists H^2$  over some half space in  $\mathbb{C}$  called  $\Omega$  and some boundary  $\partial \Omega$  in  $\mathbb{C}$ . We have the following for the functions u, v as  $u \to u|_{\partial\Omega}$  and  $v \to v|_{\partial\Omega}$ 

$$H^2(\Omega) \to H^{\frac{3}{2}}(\partial \Omega).$$

 $u \in \mathbb{C}^{\infty}(\Omega) \cap H^2(\Omega)$  and similar for  $v \in \mathbb{C}^{\infty}(\Omega) \cap H^2(\Omega)$ . This mapping extends to a unique continuous linear operator. Hence, it is onto.

 $\Rightarrow \alpha = \frac{3}{2}$  for functions or the fractional Laplace Equations for u, v exists, and moreover is Regular.

 $\Rightarrow$  smooth u, v functions. How smooth? How strong? Recall that if a function u, v or f such as  $f(z_0) \in \mathbb{C}$  domain is at least smooth enough and it meets Cauchy-Riemann conditions, then  $D^{\alpha}f(z)$  at least exists  $\forall \alpha \in \mathbb{Q}$  s.t.  $1 \leq \alpha < \infty$ . Well,  $\alpha = \frac{3}{2} \Rightarrow$  existence and more. How strong? C-R  $\Rightarrow$  strong. We use the Sharp Trace Theorem again. This time  $f \to f|_{\partial\Omega}$  and  $\exists\Omega$  and  $\partial\Omega$  in  $\mathbb{C}$ .

$$H^{\frac{3}{2}}(\Omega) \to H'(\partial\Omega)$$

 $f \subset C^{\infty}(\Omega) \cap H^{\frac{3}{2}}(\Omega)$ . This mapping is continuous linear operator. Hence, it is onto.  $\Rightarrow \alpha = 1$  which meets C-R Equations condition.

Hence,  $\alpha = \frac{3}{2}$  is strong as  $\alpha = 1$  is clearly strong. Using similar Sharp-Trace Theorem analysis on other non-integer Rationals in  $\alpha$  for  $\alpha : 1 \le \alpha \le 2$  we would see other rational order  $\alpha$  for functions f or u, v, to also be strong or at least strong enough.

Applications include the Complex or Real Functional Advection Differential Equations or FADE.

#### Example 2.4.

$$\frac{\partial C(x,t)}{\partial t} = -v \frac{\partial C(x,t)}{\partial x} + Q \frac{\partial^{\alpha} C(x,t)}{\partial x^{\alpha}}$$

 $\alpha$  is the stability or indicator of turbulence. Fourier Transforms can be used to solve.

### 3. Appendix

**Theorem 3.1** (Sharp Trace Theorem).  $\exists a \text{ half space } \Omega \text{ in } \mathbb{R}^n$ .  $\exists a \text{ boundary of a half space } \partial\Omega, \text{ also in } \mathbb{R}^n$ . Let a function  $u \to u_{\partial\Omega}$  mapping exists. Using Sobolev space  $H(\Omega)$ : we set

$$H'(\Omega) \to H^{\frac{1}{2}}(\partial\Omega)$$

This mapping for  $u \to u|_{\partial\Omega}$  extends a unique continuous linear operator so that it shows the function u is onto.  $u \in C^{\infty}(\Omega) \cap H'(\Omega)$ .

*Proof.*  $\exists$  partitions of unity  $u(x', x_n)$  s.t. defined  $x_n > 0$ 



Let  $u \in C^{\infty}(\overline{\Omega}) \cap H'(\overline{\Omega})$  be dense.

$$\hat{u}(\xi', x_n) = c \int_{\mathbb{R}^n} e^{i\xi' x'} u(x', x_n) dx', \text{(FT to } x' \text{ only)}$$

125

$$\frac{d}{dx_n}|\hat{u}|^2 = 2Re\left(\hat{u}\frac{\partial\hat{u}}{\partial x_n}\right)$$

then integrate over / from 0 to  $\infty$  or  $\int_0^\infty$ 

$$\Rightarrow = -|\hat{u}(\xi',0)|^2 = 2Re \int_0^\infty \hat{u} \frac{d\hat{u}}{dx_n} dx_n$$
$$|\hat{u}(\xi',0)|^2 \le c \int_0^\infty A|\hat{u}|^2 + \left|\frac{\partial\hat{u}}{\partial x_n}\right|^2 \frac{1}{A} dx_n$$

Now, choose  $A = [1 + |\xi'|^2]^{\frac{1}{2}} = \langle \xi' \rangle$ .

$$\therefore \langle \xi' \rangle^{\frac{1}{2}} |\hat{u}(\xi',0)|^2 \le \int_0^\infty \langle K\xi' \rangle |\hat{u}|^2 + \left| \frac{\partial \hat{u}}{\partial x_n} \right|^2 dx_n \tag{8}$$

Also, integrate (8) over  $\xi', \, \xi' \in \mathbb{R}^{n-1}$ 

$$\Rightarrow \left\| u(\cdot,0) \right\|_{H^{\frac{1}{2}}}^{2} \leq \int_{0}^{\infty} \int_{0}^{\infty} \left| (1-\Delta_{x})^{\frac{1}{2}} u \right|^{2} + \left| \frac{\partial u^{2}}{\partial x_{n}} \right| dx_{n}.$$

This is the  $\|$  trace of  $u\|_{H^{\frac{1}{2}}(\partial\Omega)}$  or

$$\begin{aligned} \|u_T\|_{H^{\frac{1}{2}}(\partial\Omega)} &\leq C \|u\|_{H'\Omega}^2 \\ \Rightarrow \text{ onto } \to \text{ by for } H'(\Omega) \to H^{\frac{1}{2}}(\partial\Omega) \end{aligned}$$

where  $T: R \to Y$ , R is in x, and dense(ness) is

$$(1 - \Delta_x)u = (1 - \Delta^5)u \Rightarrow (1 - \Delta^5)u = (1 + |\xi|^2)^u.$$

Next, use Cauchy Sequence. Prove using  $\{Tn_k\}$  converges in Y, where we define  $tn = \lim T_{n_k}$ .

Next we prove the onto of function m mapping f. Let  $f \in H^{\frac{1}{2}}(\partial\Omega)$ . Define  $v(\xi',\xi_n) = v(\xi) = \hat{f}(\xi') \frac{\langle \xi' \rangle}{\langle \xi \rangle}$ .

$$\Rightarrow \langle |\xi'| \rangle v(\xi) = \langle \xi' \rangle^{\frac{1}{2}} \cdot \hat{f}(\xi') = \frac{\langle \xi' \rangle^{\frac{1}{2}}}{\langle \xi \rangle}$$

$$\Rightarrow \langle \xi \rangle^{2} |v(\xi)|^{2} \le \left| \langle \xi' \rangle^{\frac{1}{2}} \hat{f}(\xi) \right|^{2} \cdot \frac{\langle \xi' \rangle}{\langle \xi \rangle^{2}}$$

$$\Rightarrow \int |\langle \xi \rangle| \cdot |v(\xi)|^{2} = \int_{\mathbb{R}^{n}} |\langle \xi' \rangle^{\frac{1}{2}} \hat{f}(\xi')|^{2} \otimes \left( \int_{-\infty}^{a} \frac{\langle \xi' \rangle}{\langle \xi' \rangle + \xi_{n}^{2}} d\xi_{n} d\xi' \right)$$

$$\int |\langle \xi \rangle v(\xi)|^{2} d\xi = \pi \int |\langle \xi' \rangle^{\frac{1}{2}} f|^{2} d\xi'$$

$$(**) = \pi \|f\|_{H^{\frac{1}{2}}(\partial\Omega)}^{2}$$

Let u(x) be a function s.t.  $u \in \mathbb{R}^n$ . The on a FT (Fourier Transform) is continuous

$$\Rightarrow \|u\|_{H'(\Omega)} \le c \|f\|_{H^{\frac{1}{2}}(\partial\Omega)} < \infty$$

where as an example:

$$\hat{u}(\xi',0) = Const.\hat{f}(\xi')$$

 $Lx_n = 0$ 

**Remark 3.2** (on (\*\*)).

$$\pi = \int_{-\infty}^{a} \frac{\langle \xi' \rangle}{\langle \xi' \rangle + \xi_n^2} d\xi_n$$

**Theorem 3.3** (General Trace Theorem). Let  $\Omega$  be open and the boundary  $\partial \Omega \in \mathbb{C}^{\infty}$  be continuous, s.t.

$$\left(\frac{\partial}{\partial x}\right)_{\partial\Omega}^{j}:H^{k,p}(\Omega)\to H^{l,p}(\partial\Omega)$$

where  $l = k - j - \frac{1}{p} > 0$  and p: 1 . This mapping is continuous and onto.

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