

## Solution of Non-linear Time-fractional Generalized Hirota-Satsuma Coupled Korteweg-de Vries Equation By Using New Analytical Approach

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### Abstract

This paper is concerned with the approximate analytical solution of non-linear time-fractional generalized Hirota-Satsuma Coupled Korteweg-de Vries equation (GHS-cKdV) using an efficient analytical approach, namely the Sumudu transform iterative approach. The proposed approach is an elegant amalgam of the Sumudu transform method and the Iterative method. The time-fractional derivative are described in Caputo sense. The results obtained are graphically shown and demonstrate that the approach is simple to apply and highly efficient to analyze the behavior of non-linear coupled fractional differential equations.

**Keywords:** Generalized Hirota-Satsuma coupled KdV equation; Caputo fractional derivative; Sumudu transform; Sumudu transform iterative method; Fractional differential equations.

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### 1. Introduction

Fractional calculus is a branch of mathematical study concerning derivatives and integrals of arbitrary orders. In recent years, it has been found very helpful in several fields such as rheology, viscoelasticity, biomedicine, diffusion, statistics, engineering and other fields of science ([2, 11, 21]). It is always challenging to find the exact and approximate solutions of the fractional differential equations(FDEs). Therefore it is needed to find some efficient and effective methods to solve FDEs. As a result, many different methods and approaches have been adopted so far. The Korteweg-de Vries (KdV) equations are the important equations which plays great role in different fields of science and engineering. It occurs in the study of nonlinear dispersive waves. This equation was introduced by Kortweg and de Vries for modelling of shallow water waves in canal in 1895. Firstly, Gardner *et al.* derived an analytical approach to solve this equation, which describes interactions of two long waves with

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different dispersion relations. Later, Hirota and Satsuma proposed a coupled Korteweg-de Vries (cKdV) equation

$$u_t - a(u_{\xi\xi\xi} + 6uu_{\xi}) = 2b\phi\phi_{\xi}, \quad (1)$$

$$\phi_t + \phi_{\xi\xi\xi} + 3u\phi_{\xi} = 0, \quad (2)$$

where  $a$  and  $b$  are arbitrary constants ([15, 16, 24]). In equation (1)  $2b\phi\phi_{\xi}$  acts as a force term on the KdV wave system with the linear dispersion relation  $\omega = ak^3$ . In absence of the effect of  $\phi$ , equation (1) reduced to the ordinary KdV equation

$$u_t - 6uu_{\xi} + u_{\xi\xi} = 0. \quad (3)$$

with soliton solution

$$u = 2P^2 \operatorname{sech}^2 P(x + 4aP^2 t). \quad (4)$$

Consider the Generalized Hirota-Satsuma Coupled Korteweg-de Vries (GHS-cKdV) equations

$$\begin{aligned} u_t &= \frac{1}{2}u_{\xi\xi\xi} - 3uu_{\xi} + 3(vw)_{\xi}, \\ v_t &= -v_{\xi\xi\xi} + 3uv_{\xi}, \\ w_t &= -w_{\xi\xi\xi} + 3uw_{\xi}. \end{aligned}$$

which was introduced by Wu *et al.* [22]. Recently, using an extended tanh-function method, Fan [6] constructed different soliton solutions for GHS-cKdV. In this study, we solve the GHS-cKdV of time-fractional order

$$\left. \begin{aligned} \frac{\partial^{\alpha} u}{\partial t^{\alpha}} &= \frac{1}{2} \frac{\partial^3 u}{\partial \xi^3} - 3u \frac{\partial u}{\partial \xi} + 3v \frac{\partial w}{\partial \xi} + 3w \frac{\partial v}{\partial \xi}, \\ \frac{\partial^{\alpha} v}{\partial t^{\alpha}} &= - \frac{\partial^3 v}{\partial \xi^3} + 3u \frac{\partial v}{\partial \xi}, \\ \frac{\partial^{\alpha} w}{\partial t^{\alpha}} &= - \frac{\partial^3 w}{\partial \xi^3} + 3u \frac{\partial w}{\partial \xi}, \\ 0 < \alpha < 1, \end{aligned} \right\} \quad (5)$$

subject to initial conditions

$$\left. \begin{aligned} u(\xi, 0) &= \frac{\beta - 2k^2}{3} + 2k^2 \tanh^2(k\xi), \\ v(\xi, 0) &= \frac{-4k^2 c_0 (\beta + k^2)}{3c_1^2} + \frac{4k^2 (\beta + k^2) \tanh(k\xi)}{3c_1}, \\ w(\xi, 0) &= c_0 + c_1 \tanh(k\xi), \end{aligned} \right\} \quad (6)$$

where  $k, c_0, c_1 \neq 0$  and  $\beta$  are the arbitrary constants. A lot of analytical or numerical methods have been developed to solve generalised cKdV equations such as Adomian Decomposition Method (ADM)

[4], Variational Iteration Method (VIM) [10], Reduced Differential Transform Method (RDTM) [14], Homotopy Analysis Method (HAM) [17], Homotopy Perturbation Method (HPM) [3] and many others ([5, 9, 23]). Recently, Wang and Liu used Sumudu transform (ST) in conjunction with the iterative method and became a well-known approach known as the Sumudu transform iterative method (STIM) to find approximate analytical solutions for time-fractional Cauchy reaction-diffusion equations. More recent, STIM approach was successfully implemented by Kumar and Daftardar- Gejji [13] to obtain solutions for various time and space fractional partial differential equations as well as their systems. The main objective of this study is to extend the application of Sumudu transform iterative method (STIM) to derive the approximate analytical solution of GHS-cKdV.

## 2. Preliminaries and Notations

In this portion, some fundamental definitions and properties related to fractional calculus and sumudu transform are given which are used further.

**Definition 2.1.** A fractional integral of Riemann-Liouville's order  $\alpha > 0$  of the real value function  $w(\xi, t)$  is defined as [8]

$$I_t^\alpha w(\xi, t) = \frac{1}{\Gamma(\alpha)} \int_0^t (t-s)^{\alpha-1} w(\xi, s) ds, \quad (7)$$

where  $\Gamma(\cdot)$  is known as the Gamma function.

**Definition 2.2.** The Caputo fractional derivative of function  $w(\xi, t)$  of order  $\alpha$  is defined as ([12, 18])

$$\begin{aligned} \frac{\partial^\alpha}{\partial t^\alpha} w(\xi, t) &= I_t^{m-\alpha} \left[ \frac{\partial^m}{\partial t^m} w(\xi, t) \right] \\ &= \begin{cases} \frac{1}{\Gamma(m-\alpha)} \int_0^t \frac{w^m(\xi, \tau)}{(\xi-\tau)^{\alpha-m+1}} d\tau, & m-1 < \alpha \leq m, m \in \mathbb{N}, \\ \frac{\partial^m}{\partial t^m} w(\xi, t), & \alpha = m \end{cases} \end{aligned} \quad (8)$$

**Definition 2.3.** The Sumudu transform over the set of functions  $A = \{f(t) | \exists M, \rho_1 > 0, \rho_2 > 0 \text{ such that } |f(t)| < M e^{|t|/\rho_i} \text{ if } t \in (-1)^j \times [0, \infty), j = 1, 2\}$  is defined as [7]

$$S[f(t)] = F(\omega) = \int_0^\infty e^{-t} f(\omega t) dt, \quad \omega \in (-\rho_1, \rho_2). \quad (9)$$

One of the basic property of Sumudu transform is given by

$$S\left[\frac{t^\alpha}{\Gamma(\alpha+1)}\right] = \omega^\alpha, \quad \alpha > -1. \quad (10)$$

The inverse Sumudu transform of  $\omega^\alpha$  is defined as

$$S^{-1}[\omega^\alpha] = \frac{t^\alpha}{\Gamma(\alpha+1)}, \quad \alpha > -1. \quad (11)$$

**Definition 2.4.** The Sumudu transform of Caputo time fractional derivative of  $w(\xi, t)$  of order  $\alpha > 0$  is defined as [20]

$$S\left[\frac{\partial^\alpha w(\xi, t)}{\partial t^\alpha}\right] = \omega^{-\alpha} S[w(\xi, t)] - \sum_{k=0}^{m-1} \left[ \omega^{-\alpha+k} \frac{\partial^k w(\xi, 0)}{\partial t^k} \right], \quad m-1 < \alpha \leq m, \quad m \in \mathbb{N}. \quad (12)$$

### 3. Basic Concept of STIM

In this section, the basic idea of Sumudu transform iterative method (STIM) to solve system of time and space fractional PDEs has been given. Consider the following system of general time and space fractional differential equations

$$\begin{aligned} \frac{\partial^{\gamma_i} v_i(\xi, t)}{\partial t^{\gamma_i}} &= \zeta_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right), \quad m_i - 1 < \gamma_i \leq m_i, \\ i &= 1, 2, \dots, q, \quad n-1 < \beta \leq n, \quad m_i, l, n, q \in \mathbb{N}, \end{aligned} \quad (13)$$

with initial conditions

$$\frac{\partial^j v_i(\xi, 0)}{\partial t^j} = g_{ij}(\xi), \quad j = 0, 1, 2, \dots, m_i - 1, \quad (14)$$

where  $\bar{v} = (v_1, v_2, \dots, v_q)$  and  $\zeta_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right)$  is a linear/ nonlinear operator. Taking the Sumudu transform on both sides of equation (13) and using equation (14), we have

$$S[v_i(\xi, t)] = \sum_{j=0}^{m_i-1} [w^j g_{ij}(\xi)] + w^{\gamma_i} S \left[ \zeta_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right) \right]. \quad (15)$$

On taking inverse Sumudu transform of equation (15), we get following system of equations

$$v_i(\xi, t) = S^{-1} \left( \sum_{j=0}^{m_i-1} [w^j g_{ij}(\xi)] \right) + S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right) \right) \right], \quad i = 1, 2, \dots, q. \quad (16)$$

Equation (16) can be written in following form

$$v_i(\xi, t) = f_i(\xi, t) + M_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right), \quad (17)$$

where

$$\begin{aligned} f_i(\xi, t) &= S^{-1} \left( \sum_{j=0}^{m_i-1} [w^j g_{ij}(\xi)] \right), \\ M_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right) &= S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \xi, \bar{v}, \frac{\partial^\beta \bar{v}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}}{\partial \xi^{l\beta}} \right) \right) \right]. \end{aligned} \quad (18)$$

Here  $f_i$  is known function and  $M_i$  is a linear/nonlinear operator. Equation (17) can be solved by the DJM decomposition technique introduced by Daftardar-Gejji and Jafari [19], which represents the

solution as an infinite series

$$v_i = \sum_{j=0}^{\infty} v_i^{(j)}, \quad 1 \leq i \leq q, \quad (19)$$

where the terms  $v_i^{(j)}$  are calculated recursively. Here, following abbreviations are used

$$\begin{aligned} \bar{v}^{(j)} &= (v_1^{(j)}, v_2^{(j)}, \dots, v_q^{(j)}), \\ \sum_{j=0}^r \bar{v}^{(j)} &= \left( \sum_{j=0}^r v_1^{(j)}, \sum_{j=0}^r v_2^{(j)}, \dots, \sum_{j=0}^r v_q^{(j)} \right), \quad r \in \mathbb{N} \cup \{\infty\}, \\ \frac{\partial^{k\beta} \left( \sum_{j=0}^r \bar{v}^{(j)} \right)}{\partial \xi^{k\beta}} &= \left( \frac{\partial^{k\beta} \left( \sum_{j=0}^r v_1^{(j)} \right)}{\partial \xi^{k\beta}}, \frac{\partial^{k\beta} \left( \sum_{j=0}^r v_2^{(j)} \right)}{\partial \xi^{k\beta}}, \dots, \frac{\partial^{k\beta} \left( \sum_{j=0}^r v_q^{(j)} \right)}{\partial \xi^{k\beta}} \right), \quad k \in \mathbb{N}. \end{aligned}$$

The operator  $M_i$  can be decomposed as

$$\begin{aligned} M_i \left( \zeta, \sum_{j=0}^{\infty} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{\infty} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{\infty} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) &= M_i \left( \zeta, \bar{v}^{(0)}, \frac{\partial^\beta \bar{v}^{(0)}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}^{(0)}}{\partial \xi^{l\beta}} \right) \\ &\quad + \sum_{p=1}^{\infty} \left( M_i \left( \zeta, \sum_{j=0}^p \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \\ &\quad - \sum_{p=1}^{\infty} \left( M_i \left( \zeta, \sum_{j=0}^{p-1} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right). \end{aligned} \quad (20)$$

Therefore

$$\begin{aligned} S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \zeta, \sum_{j=0}^{\infty} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{\infty} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{\infty} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \right] &= S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \zeta, \bar{v}^{(0)}, \frac{\partial^\beta \bar{v}^{(0)}}{\partial \xi^\beta}, \dots, \right. \right. \right. \\ &\quad \left. \left. \left. \frac{\partial^{l\beta} \bar{v}^{(0)}}{\partial \xi^{l\beta}} \right) \right) \right] + \sum_{p=1}^{\infty} S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \zeta, \sum_{j=0}^p \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \right] \\ &\quad - \sum_{p=1}^{\infty} S^{-1} \left[ w^{\gamma_i} S \left( \zeta_i \left( \zeta, \sum_{j=0}^{p-1} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \right]. \end{aligned} \quad (21)$$

Using equations (19) and (21) in equation (17), we get

$$\begin{aligned} \sum_{j=0}^{\infty} v_i^{(j)} &= S^{-1} \left( \sum_{j=0}^{m_i-1} [w^j g_{ij}(\xi)] \right) + S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \bar{v}^{(0)}, \frac{\partial^\beta \bar{v}^{(0)}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}^{(0)}}{\partial \xi^{l\beta}} \right) \right) \\ &+ \sum_{p=1}^{\infty} \left( S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \sum_{j=0}^p \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^p \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \right. \\ &\quad \left. - S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \sum_{j=0}^{p-1} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{p-1} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \right). \end{aligned} \quad (22)$$

The recurrence relation are defined as follows

$$\left. \begin{aligned} v_i^{(0)} &= S^{-1} \left( \sum_{j=0}^{m_i-1} [w^j g_{ij}(\xi)] \right), \\ v_i^{(1)} &= S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \bar{v}^{(0)}, \frac{\partial^\beta \bar{v}^{(0)}}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \bar{v}^{(0)}}{\partial \xi^{l\beta}} \right) \right), \\ v_i^{(m+1)} &= S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \sum_{j=0}^m \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^m \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^m \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right) \\ &- S^{-1} w^{\gamma_i} S \left( \zeta_i \left( \xi, \sum_{j=0}^{m-1} \bar{v}^{(j)}, \frac{\partial^\beta \left( \sum_{j=0}^{m-1} \bar{v}^{(j)} \right)}{\partial \xi^\beta}, \dots, \frac{\partial^{l\beta} \left( \sum_{j=0}^{m-1} \bar{v}^{(j)} \right)}{\partial \xi^{l\beta}} \right) \right), m \geq 1. \end{aligned} \right\} \quad (23)$$

Therefore, the  $m$ -term approximate analytical solution of equations (13) and (14) is given by

$$v_i \approx v_i^{(0)} + v_i^{(1)} + \dots + v_i^{(m-1)} \quad \text{or} \quad v_i \approx v_{i0} + v_{i1} + \dots + v_{i(m-1)}.$$

#### 4. The STIM Solution of the GHS-cKdV

Considering system (5), we use the application of Sumudu transform iterative method (STIM) for solving this system with initial condition (6). Now, taking the sumudu transform on the both sides of the equation (5), we obtain

$$\left. \begin{aligned} S \left[ \frac{\partial^\alpha u}{\partial t^\alpha} \right] &= S \left[ \frac{1}{2} \frac{\partial^3 u}{\partial \xi^3} - 3u \frac{\partial u}{\partial \xi} + 3v \frac{\partial w}{\partial \xi} + 3w \frac{\partial v}{\partial \xi} \right], \\ S \left[ \frac{\partial^\alpha v}{\partial t^\alpha} \right] &= S \left[ -\frac{\partial^3 v}{\partial \xi^3} + 3u \frac{\partial v}{\partial \xi} \right], \\ S \left[ \frac{\partial^\alpha w}{\partial t^\alpha} \right] &= S \left[ -\frac{\partial^3 w}{\partial \xi^3} + 3u \frac{\partial w}{\partial \xi} \right], \end{aligned} \right\} \quad (24)$$

Using formula (12), we get

$$\left. \begin{aligned} p^{-\alpha} S[u(\xi, t)] - p^{-\alpha} u(\xi, 0) &= S \left[ \frac{1}{2} \frac{\partial^3 u}{\partial \xi^3} - 3u \frac{\partial u}{\partial \xi} + 3v \frac{\partial w}{\partial \xi} + 3w \frac{\partial v}{\partial \xi} \right], \\ p^{-\alpha} S[v(\xi, t)] - p^{-\alpha} v(\xi, 0) &= S \left[ -\frac{\partial^3 v}{\partial \xi^3} + 3u \frac{\partial v}{\partial \xi} \right], \\ p^{-\alpha} S[w(\xi, t)] - p^{-\alpha} w(\xi, 0) &= S \left[ -\frac{\partial^3 w}{\partial \xi^3} + 3u \frac{\partial w}{\partial \xi} \right], \end{aligned} \right\} \quad (25)$$

so the equations become

$$\left. \begin{aligned} S[u(\xi, t)] &= u(\xi, 0) + p^\alpha S \left[ \frac{1}{2} \frac{\partial^3 u}{\partial \xi^3} - 3u \frac{\partial u}{\partial \xi} + 3v \frac{\partial w}{\partial \xi} + 3w \frac{\partial v}{\partial \xi} \right], \\ S[v(\xi, t)] &= v(\xi, 0) + p^\alpha S \left[ -\frac{\partial^3 v}{\partial \xi^3} + 3u \frac{\partial v}{\partial \xi} \right], \\ S[w(\xi, t)] &= w(\xi, 0) + p^\alpha S \left[ -\frac{\partial^3 w}{\partial \xi^3} + 3u \frac{\partial w}{\partial \xi} \right]. \end{aligned} \right\} \quad (26)$$

Taking the inverse sumudu transform of above equations, we obtain

$$\left. \begin{aligned} u(\xi, t) &= S^{-1}[u(\xi, 0)] + S^{-1} \left[ p^\alpha S \left( \frac{1}{2} \frac{\partial^3 u}{\partial \xi^3} - 3u \frac{\partial u}{\partial \xi} + 3v \frac{\partial w}{\partial \xi} + 3w \frac{\partial v}{\partial \xi} \right) \right], \\ v(\xi, t) &= S^{-1}[v(\xi, 0)] + S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3 v}{\partial \xi^3} + 3u \frac{\partial v}{\partial \xi} \right) \right], \\ w(\xi, t) &= S^{-1}[w(\xi, 0)] + S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3 w}{\partial \xi^3} + 3u \frac{\partial w}{\partial \xi} \right) \right]. \end{aligned} \right\} \quad (27)$$

In the view of recurrence relations

$$u_0 = S^{-1}[u(\xi, 0)] = \frac{\beta - 2k^2}{3} + 2k^2 \tanh^2(k\xi), \quad (28)$$

$$v_0 = S^{-1}[v(\xi, 0)] = -\frac{4k^2 c_0 (\beta + k^2)}{3c_1^2} + \frac{4k^2 (\beta + k^2) \tanh(k\xi)}{3c_1}, \quad (29)$$

$$w_0 = S^{-1}[w(\xi, 0)] = c_0 + c_1 \tanh(k\xi), \quad (30)$$

Therefore we determine the other components of the STIM solutions as follows:

First order solutions

$$\begin{aligned} u_1 &= S^{-1} p^\alpha S \left( \frac{1}{2} \frac{\partial^3 u_0}{\partial \xi^3} - 3u_0 \frac{\partial u_0}{\partial \xi} + 3v_0 \frac{\partial w_0}{\partial \xi} + 3w_0 \frac{\partial v_0}{\partial \xi} \right) \\ &= S^{-1} p^\alpha S \left( -16k^5 \tanh(k\xi) \operatorname{sech}^4(k\xi) + 8k^5 \tanh^3(k\xi) \operatorname{sech}^2(k\xi) \right. \\ &\quad \left. - 4(\beta - 2k^2)k^3 \tanh(k\xi) \operatorname{sech}^2(k\xi) - 24k^5 \tanh^3(k\xi) \operatorname{sech}^2(k\xi) \right. \\ &\quad \left. - 4k^3 \frac{c_0}{c_1} (\beta + k^2) \operatorname{sech}^2(k\xi) + 4k^3 (\beta + k^2) \tanh(k\xi) \operatorname{sech}^2(k\xi) \right. \\ &\quad \left. + (c_0 + c_1 \tanh(k\xi)) \left( \frac{4k^3 (\beta + k^2)}{c_1} \operatorname{sech}^2(k\xi) \right) \right) \end{aligned}$$

$$= 4\beta k^3 \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \quad (31)$$

$$\begin{aligned} v_1 &= S^{-1} p^\alpha S \left( -\frac{\partial^3 v_0}{\partial \xi^3} + 3u_0 \frac{\partial v_0}{\partial \xi} \right), \\ &= S^{-1} p^\alpha S \left( 8k^5 \frac{(\beta+k^2)}{3c_1} \operatorname{sech}^2(k\xi) + 4k^3(\beta-2k^2) \frac{(\beta+k^2)}{3c_1} \operatorname{sech}^2(k\xi) \right) \\ &= 4\beta k^3 \frac{(\beta+k^2)}{3c_1} \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \end{aligned} \quad (32)$$

$$\begin{aligned} w_1 &= S^{-1} p^\alpha S \left( -\frac{\partial^3 w_0}{\partial \xi^3} + 3u_0 \frac{\partial w_0}{\partial \xi} \right) \\ &= S^{-1} p^\alpha S \left( -4c_1 k^3 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) + 2c_1 k^3 \operatorname{sech}^4(k\xi) + c_1 k(\beta-2k^2) \operatorname{sech}^2(k\xi) \right. \\ &\quad \left. + 6k^3 c_1 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) \right) \\ &= c_1 \beta k \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)}. \end{aligned} \quad (33)$$

Second order solutions

$$\begin{aligned} u_2 &= S^{-1} p^\alpha S \left( \frac{1}{2} \frac{\partial^3(u_0+u_1)}{\partial \xi^3} - 3(u_0+u_1) \frac{\partial(u_0+u_1)}{\partial \xi} + 3(v_0+v_1) \frac{\partial(w_0+w_1)}{\partial \xi} \right. \\ &\quad \left. + 3(w_0+w_1) \frac{\partial(v_0+v_1)}{\partial \xi} \right) - S^{-1} p^\alpha S \left( \frac{1}{2} \frac{\partial^3 u_0}{\partial \xi^3} - 3u_0 \frac{\partial u_0}{\partial \xi} + 3v_0 \frac{\partial w_0}{\partial \xi} + 3w_0 \frac{\partial v_0}{\partial \xi} \right) \\ &= S^{-1} p^\alpha S \left( 4k^3(\beta+k^2) \operatorname{sech}^2(k\xi) \tanh(k\xi) + 4\beta^2 k^4 \operatorname{sech}^4(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} - 8\beta^2 k^4 \operatorname{sech}^2(k\xi) \right. \\ &\quad \left. \tanh^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} - 16k^5 \operatorname{sech}^2(k\xi) \tanh(k\xi) - 16\beta^3 k^5 \operatorname{sech}^4(k\xi) \tanh(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} \right. \\ &\quad \left. - 16\beta^2 k^7 \operatorname{sech}^4(k\xi) \tanh(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} + 16\beta k^6 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \right. \\ &\quad \left. - 16\beta k^6 \operatorname{sech}^6(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} + 16\beta k^6 \operatorname{sech}^4(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} + 32\beta k^6 \operatorname{sech}^2(k\xi) \tanh^4(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \right. \\ &\quad \left. - 32\beta k^6 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} + 96\beta^2 k^7 \operatorname{sech}^4(k\xi) \tanh^3(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} \right. \\ &\quad \left. - 48\beta^2 k^7 \operatorname{sech}^6(k\xi) \tanh(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} \right) - 4\beta k^3 \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \\ &= 4k^4 \beta^2 \operatorname{sech}^2(k\xi) (\operatorname{sech}^2(k\xi) - 2\tanh^2(k\xi)) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \\ &\quad - 16\beta^3 k^5 \operatorname{sech}^4(k\xi) \tanh(k\xi) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \frac{\Gamma(2\alpha+1)}{(\Gamma(\alpha+1))^2} \\ &\quad + 16\beta^2 k^7 \operatorname{sech}^4(k\xi) \tanh(k\xi) (5\tanh^2(k\xi) - 4\operatorname{sech}^2(k\xi)) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \frac{\Gamma(2\alpha+1)}{(\Gamma(\alpha+1))^2} \end{aligned} \quad (34)$$

$$\begin{aligned} v_2 &= S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3(v_0+v_1)}{\partial \xi^3} + 3(u_0+u_1) \frac{\partial(v_0+v_1)}{\partial \xi} \right) \right] - S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3 v_0}{\partial \xi^3} + 3u_0 \frac{\partial v_0}{\partial \xi} \right) \right] \\ &= S^{-1} p^\alpha S \left( \frac{4k^2(\beta+k^2)}{3c_1} [2k^3 \operatorname{sech}^4(k\xi) - 2k^3 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) + 8\beta k^4 \operatorname{sech}^2(k\xi) \tanh(k\xi) \right. \\ &\quad \left. (\tanh^2(k\xi) - 2\operatorname{sech}^2(k\xi)) \frac{t^\alpha}{\Gamma(\alpha+1)} + (\beta-2k^2) k \operatorname{sech}^2(k\xi) - 2\beta k^2 (\beta-2k^2) \operatorname{sech}^2(k\xi) \tanh(k\xi) \right. \\ &\quad \left. \frac{t^\alpha}{\Gamma(\alpha+1)} + 6k^3 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) - 12\beta k^4 \operatorname{sech}^2(k\xi) \tanh^3(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} + 12\beta k^4 \operatorname{sech}^4(k\xi) \tanh(k\xi) \right) \end{aligned}$$

$$\begin{aligned} & \frac{t^\alpha}{\Gamma(\alpha+1)} - 24\beta^2 k^5 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} \Big] - \frac{4\beta k^3(\beta+k^2)}{3c_1} \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \\ & = -\frac{8k^4\beta^2(\beta+k^2)}{3c_1} \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \\ & \quad - \frac{32}{c_1} \beta^2 k^7 (\beta+k^2) \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{3\alpha} \Gamma(2\alpha+1)}{\Gamma(3\alpha+1)(\Gamma(\alpha+1))^2} \end{aligned} \quad (35)$$

$$\begin{aligned} w_2 &= S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3(w_0+w_1)}{\partial\xi^3} + 3(u_0+u_1) \frac{\partial(w_0+w_1)}{\partial\xi} \right) \right] - S^{-1} \left[ p^\alpha S \left( -\frac{\partial^3 w_0}{\partial\xi^3} + 3u_0 \frac{\partial w_0}{\partial\xi} \right) \right] \\ &= S^{-1} \left[ p^\alpha S \left( 2c_1 k^3 \operatorname{sech}^4(k\xi) - 4c_1 k^3 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) - 8k^4 c_1 \beta \operatorname{sech}^2(k\xi) \tanh(k\xi) \right. \right. \\ &\quad (2\operatorname{sech}^2(k\xi) - \tanh^2(k\xi)) \frac{t^\alpha}{\Gamma(\alpha+1)} + (\beta - 2k^2) k c_1 \operatorname{sech}^2(k\xi) - 2c_1 \beta k^2 (\beta - 2k^2) \\ &\quad \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} + 6k^3 c_1 \operatorname{sech}^2(k\xi) \tanh^2(k\xi) + 12k^4 c_1 \beta \operatorname{sech}^2(k\xi) \tanh(k\xi) \\ &\quad (\operatorname{sech}^2(k\xi) - \tanh^2(k\xi)) \frac{t^\alpha}{\Gamma(\alpha+1)} - 24c_1 \beta^2 k^5 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{2\alpha}}{(\Gamma(\alpha+1))^2} \Big) \\ &\quad \left. - c_1 \beta k \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \right] \\ &= -2c_1 \beta^2 k^2 \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} - 24c_1 \beta^2 k^5 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{3\alpha} \Gamma(2\alpha+1)}{\Gamma(3\alpha+1)(\Gamma(\alpha+1))^2}. \end{aligned} \quad (36)$$

Similarly more iterations can be obtained. The STIM yields the solution  $u(\xi, t), v(\xi, t), w(\xi, t)$  as follows

$$\begin{aligned} u(\xi, t) &= u_0 + u_1 + u_2 + \dots \\ &= \frac{\beta - 2k^2}{3} + 2k^2 \tanh^2(k\xi) + 4\beta k^3 \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \\ &\quad + 4k^4 \beta^2 \operatorname{sech}^2(k\xi) (\operatorname{sech}^2(k\xi) - 2\tanh^2(k\xi)) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \\ &\quad - 16\beta^3 k^5 \operatorname{sech}^4(k\xi) \tanh(k\xi) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \frac{\Gamma(2\alpha+1)}{(\Gamma(\alpha+1))^2} \\ &\quad + 16\beta^2 k^7 \operatorname{sech}^4(k\xi) \tanh(k\xi) (5\tanh^2(k\xi) - 4\operatorname{sech}^2(k\xi)) \frac{t^{3\alpha}}{\Gamma(3\alpha+1)} \frac{\Gamma(2\alpha+1)}{(\Gamma(\alpha+1))^2} + \dots \end{aligned} \quad (37)$$

$$\begin{aligned} v(\xi, t) &= v_0 + v_1 + v_2 + \dots \\ &= -\frac{4k^2 c_0 (\beta + k^2)}{3c_1^2} + \frac{4k^2 (\beta + k^2) \tanh(k\xi)}{3c_1} + 4\beta k^3 \frac{(\beta + k^2)}{3c_1} \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} \\ &\quad - \frac{8k^4 \beta^2 (\beta + k^2)}{3c_1} \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \\ &\quad - \frac{32}{c_1} \beta^2 k^7 (\beta + k^2) \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{3\alpha} \Gamma(2\alpha+1)}{\Gamma(3\alpha+1)(\Gamma(\alpha+1))^2} + \dots \end{aligned} \quad (38)$$

$$\begin{aligned} w(\xi, t) &= w_0 + w_1 + w_2 + \dots \\ &= c_0 + c_1 \tanh(k\xi) + c_1 \beta k \operatorname{sech}^2(k\xi) \frac{t^\alpha}{\Gamma(\alpha+1)} - 2c_1 \beta^2 k^2 \tanh(k\xi) \operatorname{sech}^2(k\xi) \frac{t^{2\alpha}}{\Gamma(2\alpha+1)} \\ &\quad - 24c_1 \beta^2 k^5 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) \frac{t^{3\alpha} \Gamma(2\alpha+1)}{\Gamma(3\alpha+1)(\Gamma(\alpha+1))^2} + \dots \end{aligned} \quad (39)$$

**Special case: for  $\alpha = 1$**

$$u = \frac{\beta - 2k^2}{3} + 2k^2 \tanh^2(k\xi) + 4\beta k^3 \tanh(k\xi) \operatorname{sech}^2(k\xi)t + 2k^4 \beta^2 \operatorname{sech}^2(k\xi) (\operatorname{sech}^2(k\xi) - 2\tanh^2(k\xi)) t^2 - \frac{16}{3} \beta^3 k^5 \operatorname{sech}^4(k\xi) \tanh(k\xi) t^3 + \frac{16}{3} \beta^2 k^7 \operatorname{sech}^4(k\xi) \tanh(k\xi) (5\tanh^2(k\xi) - 4\operatorname{sech}^2(k\xi)) t^3 + \dots \quad (40)$$

$$v = -\frac{4k^2 c_0 (\beta + k^2)}{3c_1^2} + \frac{4k^2 (\beta + k^2)}{3c_1} \tanh(k\xi) + 4\beta k^3 \frac{(\beta + k^2)}{3c_1} \operatorname{sech}^2(k\xi)t - \frac{4k^4 \beta^2 (\beta + k^2)}{3c_1} \tanh(k\xi) \operatorname{sech}^2(k\xi)t^2 - \frac{32}{3c_1} \beta^2 k^7 (\beta + k^2) \operatorname{sech}^4(k\xi) \tanh^2(k\xi) t^3 + \dots \quad (41)$$

$$w = c_0 + c_1 \tanh(k\xi) + c_1 \beta k \operatorname{sech}^2(k\xi)t - c_1 \beta^2 k^2 \tanh(k\xi) \operatorname{sech}^2(k\xi)t^2 - 8c_1 \beta^2 k^5 \operatorname{sech}^4(k\xi) \tanh^2(k\xi) t^3 + \dots \quad (42)$$

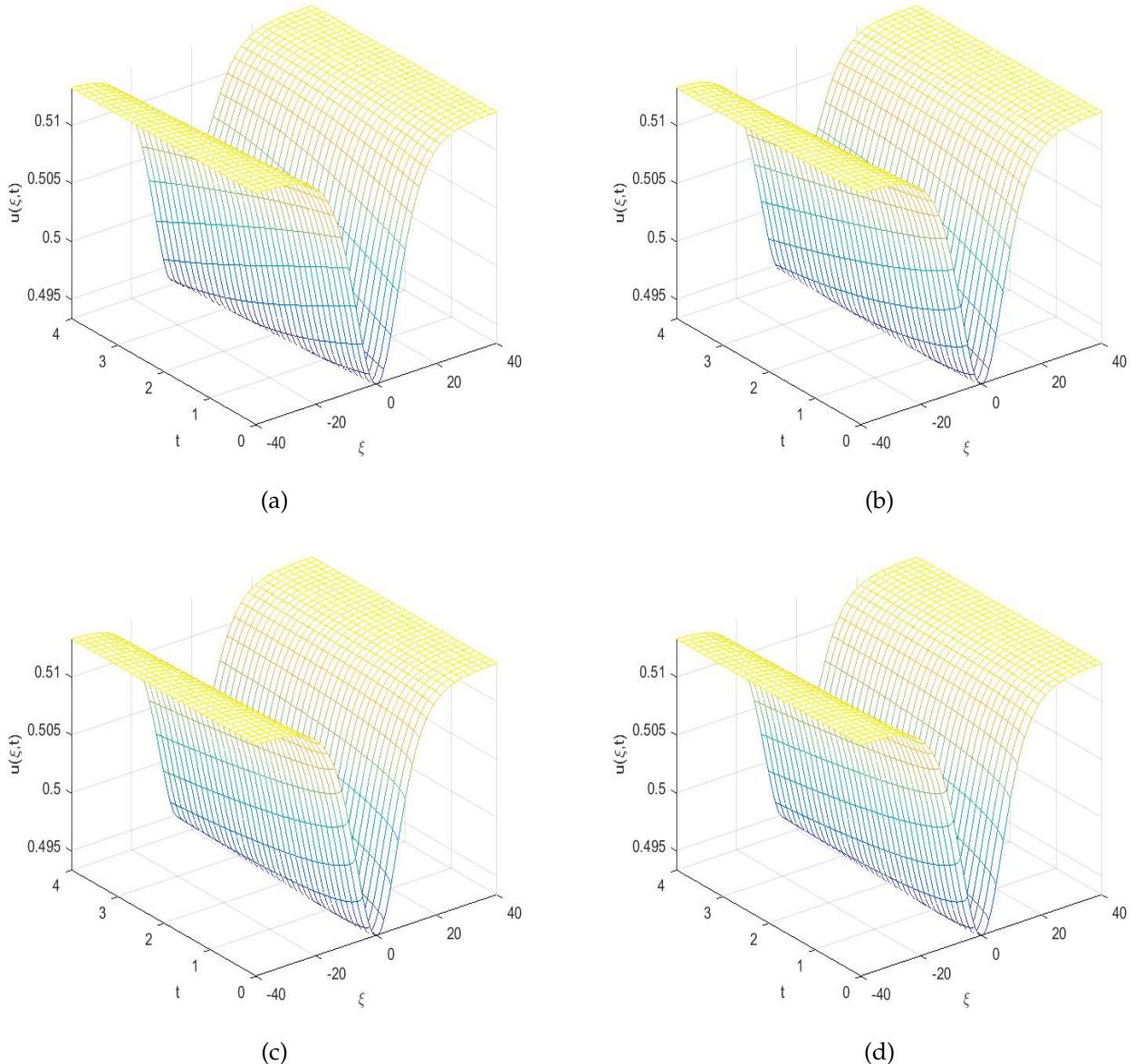


Figure 1: The surface shows the solution  $u(\xi, t)$ , when  $\beta = 1.5, k = 0.1$ , : (a) Approximate solution for  $\alpha = 1$ , (b) Approximate solution for  $\alpha = 0.75$ , (c) Approximate solution for  $\alpha = 0.50$ , (d) Approximate solution for  $\alpha = 0.25$

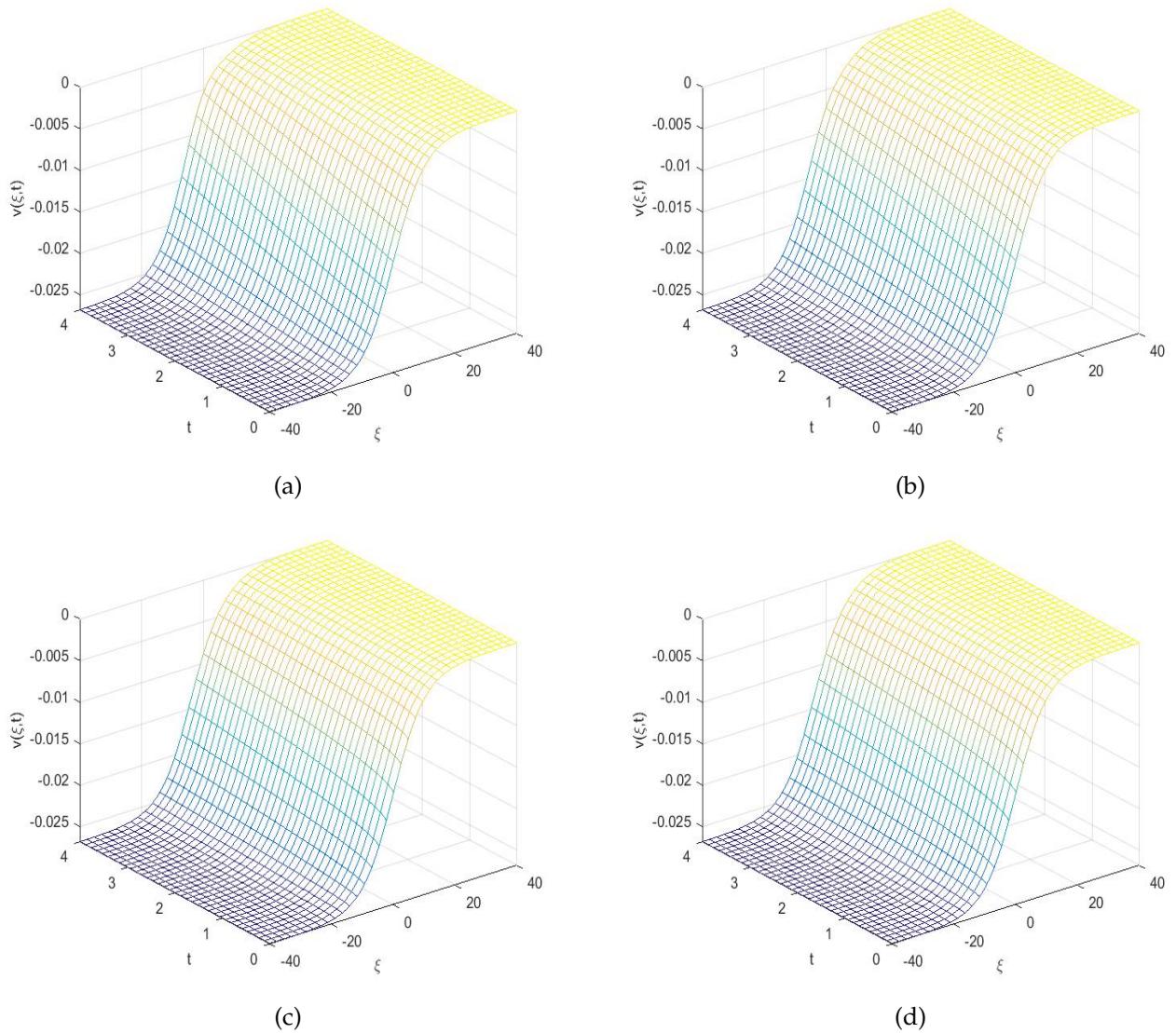
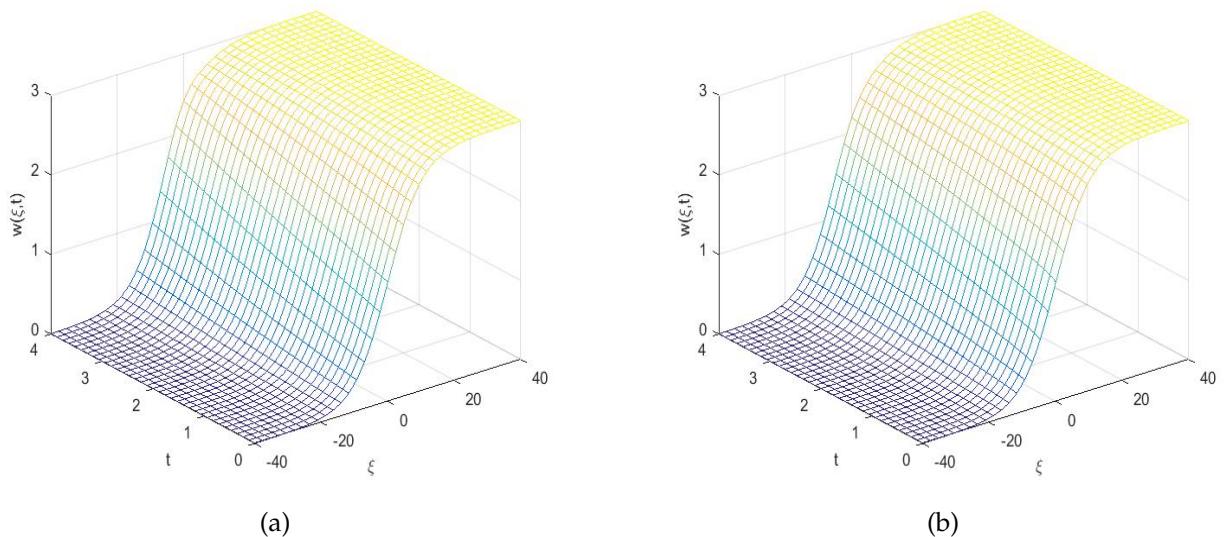


Figure 2: The surface shows the solution  $v(\xi, t)$ , when  $\beta = 1.5$ ,  $k = 0.1$ ,  $c_0 = 1.5, c_1 = 1.5$  : (a) Approximate solution for  $\alpha = 1$ , (b) Approximate solution for  $\alpha = 0.75$ , (c) Approximate solution for  $\alpha = 0.50$ , (d) Approximate solution for  $\alpha = 0.25$



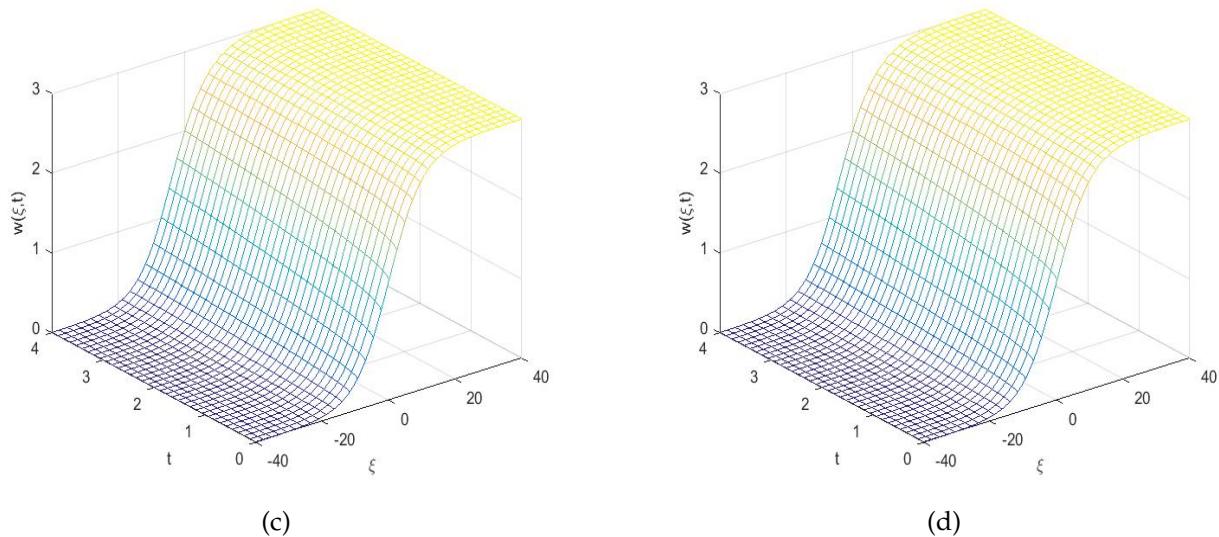


Figure 3: The surface shows the solution  $w(\xi, t)$ , when  $\beta = 1.5$ ,  $k = 0.1$ ,  $c_0 = 1.5$ ,  $c_1 = 1.5$ : (a) Approximate solution for  $\alpha = 1$ , (b) Approximate solution for  $\alpha = 0.75$ , (c) Approximate solution for  $\alpha = 0.50$ , (d) Approximate solution for  $\alpha = 0.25$

## 5. Conclusion

In this paper, the time-fractional generalised Hirota-Satsuma coupled Kortweg-de Vries (GHS-cKdV) equations were solved by new approach namely Sumudu Transform Iterative Method (STIM) within Caputo operator. By increasing the number of iteration , the accuracy of STIM can be further improved.

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